

1. This problem is about the orbital angular momentum operator, $\hat{L}_i = \epsilon_{ijk} \hat{X}_j \hat{P}_k$.

(a) Using canonical commutation relations between components of \hat{X}_i and \hat{P}_i , show that

$$\left[\hat{X}_i, \hat{L}_j \right] = i\hbar \epsilon_{ijk} \hat{X}_k \quad \text{and} \quad \left[\hat{P}_i, \hat{L}_j \right] = i\hbar \epsilon_{ijk} \hat{P}_k . \quad (1)$$

(b) Show that

$$\left[\hat{L}_i, \hat{L}_j \right] = i\hbar \epsilon_{ijk} \hat{L}_k \quad \text{and} \quad \epsilon_{ijk} \hat{L}_j \hat{L}_k = i\hbar \hat{L}_i . \quad (2)$$

(c) Define $\hat{L}_\pm = \hat{L}_x \pm i\hat{L}_y$ and show that

$$\left[\hat{L}_z, \hat{L}_\pm \right] = \pm\hbar \hat{L}_\pm \quad \text{and} \quad \left[\hat{L}_+, \hat{L}_- \right] = 2\hbar \hat{L}_z . \quad (3)$$

(d) Show that in the spherical coordinate basis

$$\begin{aligned} \hat{L}_z \Psi(r, \theta, \phi) &= -i\hbar \frac{\partial}{\partial \phi} \Psi(r, \theta, \phi) , \\ \hat{L}_\pm \Psi(r, \theta, \phi) &= -i\hbar e^{\pm i\phi} \left(\pm i \frac{\partial}{\partial \theta} - \frac{1}{\tan \theta} \frac{\partial}{\partial \phi} \right) \Psi(r, \theta, \phi) . \end{aligned} \quad (4)$$

(e) Compute $\hat{L}_i \hat{L}_i \Psi(r, \theta, \phi)$ in the same basis.

(Hint: use $\hat{L}_i \hat{L}_i = \hat{L}_z^2 + \frac{1}{2} \hat{L}_+ \hat{L}_- + \frac{1}{2} \hat{L}_- \hat{L}_+$.)

2. This problem is about rotations in three dimensions. For a rotation by angle α around an axis of a unit vector \vec{n} , the rotation matrix is

$$R_{ij}(\alpha, \vec{n}) = \delta_{ij} \cos \alpha - \epsilon_{ijk} n_k \sin \alpha + n_i n_j (1 - \cos \alpha) . \quad (5)$$

A product of two rotations $R(\alpha', \vec{n}') \cdot R(\alpha'', \vec{n}'')$ is itself a rotation by some angle α around some axis \vec{n} . To determine α and \vec{n} , we demand

$$R_{ij}(\alpha, \vec{n}) = R_{ik}(\alpha', \vec{n}') \cdot R_{jk}(\alpha'', \vec{n}'') , \quad (6)$$

and then substitute eq. (5) and solve for α and \vec{n} .

(a) Show that to second order in α' and α'' ,

$$\alpha \vec{n} \simeq \alpha' \vec{n}' + \alpha'' \vec{n}'' + \frac{1}{2} \alpha' \alpha'' \vec{n}' \times \vec{n}'' . \quad (7)$$

(b) Prove $\text{tr} R \equiv R_{ii} = 1 + 2 \cos \alpha$ and $\epsilon_{ijk} R_{jk} = -2n_i \sin \alpha$ and use these formulas to derive exact expressions for $\cos \alpha$ and $\vec{n} \sin \alpha$ in terms of α' , α'' , \vec{n}' and \vec{n}'' .

3. This problem is about Pauli matrices and the relation between abstract rotations in spin space and rotations in three dimensions.

(a) Show that for any two vectors \vec{a} and \vec{b} ,

$$(\vec{a} \cdot \vec{\sigma})(\vec{b} \cdot \vec{\sigma}) = (\vec{a} \cdot \vec{b})\mathbf{1} + i(\vec{a} \times \vec{b}) \cdot \vec{\sigma}. \quad (8)$$

(b) Show that for any angle α and for any unit vector \vec{n} ,

$$U(\alpha, \vec{n}) \equiv \exp(-i\frac{\alpha}{2}\vec{n} \cdot \vec{\sigma}) = \cos(\frac{\alpha}{2})\mathbf{1} - i\sin(\frac{\alpha}{2})\vec{n} \cdot \vec{\sigma}. \quad (9)$$

(c) Solve for α and \vec{n} such that

$$U(\alpha, \vec{n}) = U(\alpha', \vec{n}') U(\alpha'', \vec{n}''). \quad (10)$$

(d) The commutation relations between generators of a Lie algebra completely determine the products of their exponentials. Consider a vector of 2×2 hermitian matrices, $\vec{S} \equiv \frac{\hbar}{2}\vec{\sigma}$; commutation relations between its components are exactly the same as between the components \hat{J}_i of the angular momentum operator. Since the definition of the U matrices amounts to $U(\alpha, \vec{n}) = \exp(\alpha\vec{n} \cdot \vec{S}/i\hbar)$, which is precisely analogous to $\hat{R}(\alpha, \vec{n}) = \exp(\alpha\vec{n} \cdot \vec{J}/i\hbar)$, we expect the product rule for the U matrices to be exactly the same as for the unitary rotation operators \hat{R} , which in turn could be identical to the product rule for the R_{ij} rotation matrices given in eq. (6).

Verify that a solution of eq. (10) is automatically a solution of eq. (6).

Is the converse also true?