

Now insert  $TT = \mathbb{1}$  everywhere and apply rules:

$$T\psi(t, \mathbf{x})T = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s \left( a_{-\mathbf{p}}^{-s} [U^s(\mathbf{p})]^* e^{i\mathbf{p}\cdot\mathbf{x}} + b_{-\mathbf{p}}^{-s} [V^s(\mathbf{p})]^* e^{-i\mathbf{p}\cdot\mathbf{x}} \right)$$

Action of  $T$  on operators
Action of  $T$  on c-numbers

Now change variables  $\mathbf{p} \rightarrow \tilde{\mathbf{p}}$  and use  $[U]^*$ ,  $[V]^*$  relations:

$$T\psi(t, \mathbf{x})T = (-\gamma_1 \gamma_3) \int \frac{d^3\tilde{\mathbf{p}}}{(2\pi)^3} \frac{1}{\sqrt{2E_{\tilde{\mathbf{p}}}}} \sum_s \left( a_{\tilde{\mathbf{p}}}^{-s} U^{-s}(\tilde{\mathbf{p}}) e^{i\tilde{\mathbf{p}}\cdot(t, -\mathbf{x})} + b_{\tilde{\mathbf{p}}}^{-s} V^{-s}(\tilde{\mathbf{p}}) e^{-i\tilde{\mathbf{p}}\cdot(t, -\mathbf{x})} \right)$$

$$\left\{ \text{Use } \tilde{\mathbf{p}}\cdot(t, -\mathbf{x}) = (p^0, -\mathbf{p})\cdot(t, -\mathbf{x}) = -\tilde{\mathbf{p}}\cdot(-t, \mathbf{x}) \right\}$$

$$\Rightarrow T\psi(\mathbf{x}, t)T = (-\gamma_1 \gamma_3) \psi(-t, \mathbf{x})$$

Simple trace trick on  $\psi$  !!

$$T\bar{\psi}T = (T\psi T)^\dagger (\gamma^0)^* = \psi^\dagger(-t, \mathbf{x}) [-\gamma_1 \gamma_3]^\dagger \gamma^0$$

$$\Rightarrow T\bar{\psi}(\mathbf{x}, t)T = \bar{\psi}(-t, \mathbf{x}) \gamma_1 \gamma_3$$

Now check action of  $T$  on various bilinears.

SCALAR

$$\begin{aligned} T \bar{\psi} \psi(t, x_i) T &= \bar{\psi} \gamma^1 \gamma^3 (-\gamma^1 \gamma^3) \psi(-t, x_i) \\ &= \underline{+ \bar{\psi} \psi(-t, x_i)} \end{aligned}$$

PSEUDO SCALAR

$$\begin{aligned} T i \bar{\psi} \gamma^5 \psi T &= -i \bar{\psi} (\gamma^1 \gamma^3) \gamma^5 (-\gamma^1 \gamma^3) \psi(-t, x_i) \\ &= \underline{-i \bar{\psi} \gamma^5 \psi(-t, x_i)} \end{aligned}$$

VECTOR

$$\begin{aligned} T \bar{\psi} \gamma^\mu \psi T &= \bar{\psi} \gamma^1 \gamma^3 (\gamma^\mu)^\dagger (-\gamma^1 \gamma^3) \psi(-t, x_i) \\ &= \begin{cases} + \bar{\psi} \gamma^\mu \psi(-t, x_i) & \mu = 0 \\ - \bar{\psi} \gamma^\mu \psi(-t, x_i) & \mu = i \end{cases} \end{aligned}$$

AXIAL VECTOR

SAME AS VECTOR!

## Charge Conjugation

Define:

$$C a_p^s C = b_p^s, \quad C b_p^s C = a_p^s$$

Takes fermion into antifermion w/ same spin  
 $\{$  " antifermion " fermion w/ " "

What is action of C on  $\psi$ ??

Recall:

$$u^s(p) = \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix}, \quad v^s(p) = \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^{-s} \\ -\sqrt{p \cdot \bar{\sigma}} \xi^{-s} \end{pmatrix}$$

$$\xi^{-s} = -i \sigma^2 (\xi^s)^*$$

$$\sqrt{p \cdot \sigma} \sigma^2 = \sigma^2 \sqrt{p \cdot \bar{\sigma}}$$

$$\sqrt{p \cdot \bar{\sigma}} \sigma^2 = \sigma^2 \sqrt{p \cdot \sigma}$$

$$(v^s(p))^* = \begin{pmatrix} \sqrt{p \cdot \sigma} (-i \sigma^2 (\xi^s)^*) \\ -\sqrt{p \cdot \bar{\sigma}} (-i \sigma^2 (\xi^s)^*) \end{pmatrix}^*$$

$$= \begin{pmatrix} -i \sigma^2 \sqrt{p \cdot \bar{\sigma}} \xi^{s*} \\ i \sigma^2 \sqrt{p \cdot \sigma} \xi^{s*} \end{pmatrix}^* = \begin{pmatrix} -i \sigma^2 \sqrt{p \cdot \bar{\sigma}} \xi^s \\ i \sigma^2 \sqrt{p \cdot \sigma} \xi^s \end{pmatrix}$$

$$= \begin{pmatrix} 0 & -i \sigma^2 \\ i \sigma^2 & 0 \end{pmatrix} \begin{pmatrix} \sqrt{p \cdot \sigma} \xi^s \\ \sqrt{p \cdot \bar{\sigma}} \xi^s \end{pmatrix}$$

$$\Rightarrow (U^s(p_i))^* = -i\gamma^2 U^s(p_i)$$

$$\Rightarrow U^s(p) = -i\gamma^2 (U^s(p))^*$$

Similarly,

$$V^s(p) = -i\gamma^2 (U^s(p))^*$$

(simply c.c. of above!!)

we're ready!

$$C \psi(x) C = \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}}$$

$$\times \sum_s \left( -i\gamma^2 (V^s(p))^* b_r^s e^{-ip \cdot x} - i\gamma^2 (U^s(p))^* a_r^{s\dagger} e^{ip \cdot x} \right)$$

$$= -i\gamma^2 \psi^*(x) = -i\gamma^2 (\psi^\dagger)^T$$

$$\Rightarrow C \psi(x) C = -i (\bar{\psi} \gamma_0 \gamma^2)^T$$

$$C \bar{\psi}(x) C = C \psi^\dagger C \gamma_0 = (-i\gamma^2 \psi)^T \gamma_0$$

$$\Rightarrow C \bar{\psi}(x) C = (-i \gamma_0 \gamma^2 \psi)^T$$

(Note that  $C$  is a linear unitary operator even though takes  $\psi \rightarrow \psi^c$ )

Action on bilinears?

### SCALAR

$$\begin{aligned}
 C \bar{\psi} \psi C &= (-i \gamma^0 \gamma^2 \psi)^T (-i \bar{\psi} \gamma^0 \gamma^2)^T \\
 &= -\gamma_{cb}^0 \gamma_{bc}^2 \psi_c \bar{\psi}_d \gamma_{de}^0 \gamma_{ea}^2 \\
 &= +\bar{\psi}_d \gamma_{de}^0 \gamma_{ea}^2 \gamma_{cb}^0 \gamma_{bc}^2 \psi_c \\
 &= \bar{\psi} \gamma^0 \gamma^2 \gamma^0 \gamma^2 \psi = -\bar{\psi} \gamma^0 \gamma^2 \gamma^0 \gamma^2 \psi \\
 &= \underline{+\bar{\psi} \psi}
 \end{aligned}$$

Similarly one finds:

### PSEUDOSCALAR

$$C i \bar{\psi} \gamma^5 \psi C = \underline{i \bar{\psi} \gamma^5 \psi}$$

### VECTOR

$$C \bar{\psi} \gamma^\mu \psi C = \underline{-\bar{\psi} \gamma^\mu \psi}$$

### AXIAL VECTOR

$$C \bar{\psi} \gamma^\mu \gamma^5 \psi C = \underline{+\bar{\psi} \gamma^\mu \gamma^5 \psi}$$

Now let's summarize our discrete transformations!

Summary

	$\bar{\psi}\psi$	$i\bar{\psi}\gamma_5\psi$	$\bar{\psi}\gamma^4\psi$	$\bar{\psi}\gamma^4\gamma_5\psi$	$\partial_\mu$
P	+1	-1	$(-1)^n$	$-(-1)^n$	$(-1)^n$
T	+1	-1	$(-1)^n$	$(-1)^n$	$-(-1)^n$
C	+1	+1	-1	+1	+1
<hr/>					
CPT	+1	+1	-1	-1	-1

Note:  $\mathcal{L} = \bar{\psi} (i\gamma^\mu \partial_\mu - m)\psi$  is CPT invariant

CPT theorem: (Pati)

One cannot build a Lorentz invariant CFT w/ a Hermitian Hamiltonian that violates CPT.

(See Itzykson + Zuber for sketch of proof!)

Basically want to show that

$CPT \mathcal{L}(x) (CPT)^\dagger = \mathcal{L}(-x)$