

Path Integrals in Quantum Mechanics

- Useful to see alternate derivation of Feynman Rules
- Close analogy between Q.F.T. and STAT MECH.

Consider non-relativistic Q.M. in one dimension:

$$H = \frac{p^2}{2m} + V(x)$$

Suppose we wish to describe a particle travelling from position x_a to position x_b in time T .

Amplitude for this process is

$$U(x_a, x_b; T) = \langle x_b | e^{-iHT/\hbar} | x_a \rangle$$

(in "canonical" Hamiltonian formalism)

Let's motivate a different mathematical expression for this amplitude.

Superposition principle: amplitude is coherent sum of all ways in which process can take place.

In our example, consider the amplitude for each particular spatial path that particle can take as a pure phase. Then expect:

$$U(x_a, x_b; T) = \sum_{\text{all paths}} e^{i \text{phase}} = \int D[x(t)] e^{i \text{phase}}$$

indicates 1 path for every function $x(t)$ that begins at x_a and ends at x_b

Aside on Functionals

Functional: A function that maps functions to numbers.

e.g. $F[x(t)]$ is a functional.
 (note square brackets)
 Can be integrated over a set of functions $x(t)$.

We will see more properties of functionals as we proceed.

What do we use for "phase" in $\int D[x(t)]$??

In classical limit there should be 1 path: the classical path

One possibility is that "stationary phase approximation" picks out classical path.

Recall: $\mathbb{T} f \quad \mathbb{T}(k) = \int dx g(x) e^{ikf(x)}$

S.P.A. $\Rightarrow \mathbb{T}(k)$ depends only on critical points of $f(x)$:
 i.e. when $f'(x)|_{x=x_c} = 0$

Hence would expect classical path $x_{cl}(t)$ to be given by:

$$\frac{\delta}{\delta x(t)} (\text{phase}[x(t)]) \Big|_{x=x_{cl}} = 0$$

However, we also know that

$$\frac{\delta}{\delta x(t)} S[x(t)] \Big|_{x=x_{cl}} = 0$$

"principle of least action"

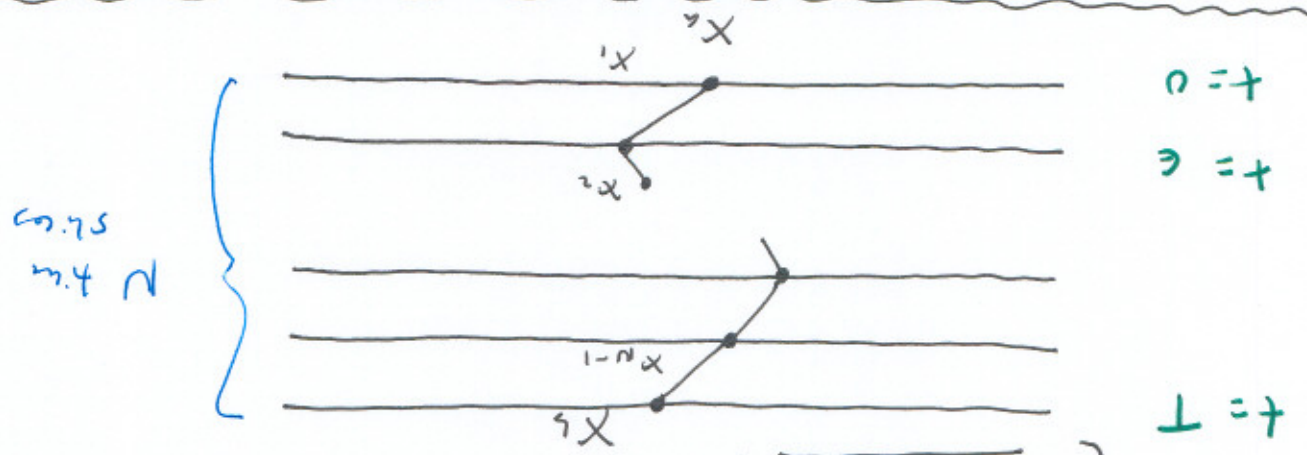
$$S = \int dt L$$

\Rightarrow $\text{phase} = \frac{S}{\hbar}$ ← constant w/ units of action

{ S.P.A. valid when $S \gg \hbar$ }

$$S = \int_T^0 dt + \left(\frac{1}{m} \dot{x}^2 - V(x) \right) - \sum_{i=1}^k \left(\frac{1}{m} \dot{x}_{i+1}^2 - \frac{e}{\hbar} V(x_{i+1} + x_k) \right)$$

Discretized action:



Use discretization: break up time interval $0 \leq T$ into many small pieces of duration ϵ .
 Approximate a "path" $x(t)$ as a sequence of straight lines, one for each thin slice.

define: $\int D x(t)$

To evaluate the path integral explicitly, we must

canonical Hamiltonian formula of Q.M.
 path integral formula of Q.M.

$$\langle X_f | e^{-iHT} | X_i \rangle = U(x_f, X_f; T) = \int D x(t) e^{iS[x(t)]/\hbar}$$

\therefore we have the conjecture:

Define the path integral measure by:

$$\int D\psi(\epsilon) \equiv \frac{1}{C(\epsilon)} \int \frac{dx_1}{C(\epsilon)} \int \frac{dx_2}{C(\epsilon)} \dots \int \frac{dx_{N-1}}{C(\epsilon)} = \frac{1}{C(\epsilon)^N} \prod_k \int_{-\infty}^{\infty} \frac{dx_k}{C(\epsilon)}$$

$C(\epsilon) = \text{constant}$

Reason for this definition will become clear below

Now we would like to show equivalence of "canonical" and "path integral" approaches for 1 particle potential problem by showing that both amplitudes satisfy the same differential equation, with some initial condition.

Consider the addition of the very last time slice in discretized sum over paths:

$$U(x_a, x_b; T) = \int_{-\infty}^{\infty} \frac{dx'}{C(\epsilon)} \exp\left[\frac{i}{\hbar} m \frac{(x_b - x')^2}{2\epsilon} - \frac{i}{\hbar} \epsilon V\left(\frac{x_b + x'}{2}\right)\right] U(x_a, x'; T - \epsilon)$$

" $U(x'; x_b; \epsilon)$

Now consider limit $\epsilon \rightarrow 0$

oscillates rapidly!!

constrain x' to be near x_b ! \Rightarrow $\epsilon \propto V(x_b)$!

Expanded in powers of $\delta \equiv (x' - x_0)$:

$$\left\{ \begin{aligned} V(x_0 + x') &= V(x_0 + \delta/2) \approx V(x_0) \\ U(x_0, x'; T-\epsilon) &= U(x_0, x_0 + \delta; T-\epsilon) \\ &\approx \left[1 + \delta \frac{\partial}{\partial x_0} + \frac{1}{2} \delta^2 \frac{\partial^2}{\partial x_0^2} + \dots \right] U(x_0, x_0; T-\epsilon) \end{aligned} \right.$$

$$U(x_0, x_0; T) = \int_{-\infty}^{\infty} \frac{dx'}{c} \exp\left(\frac{i}{\hbar} \frac{m}{2\epsilon} (x_0 - x')^2\right) \left[1 - \frac{i\epsilon}{\hbar} V(x_0) \right] \\ \times \left[1 + (x' - x_0) \frac{\partial}{\partial x_0} + \frac{1}{2} (x' - x_0)^2 \frac{\partial^2}{\partial x_0^2} + \dots \right] U(x_0, x_0; T-\epsilon)$$

We can now do the "Gaussian" integrals over x' !!

$$\left\{ \int dy e^{-by^2} = \sqrt{\frac{\pi}{b}} \quad \int dy y e^{-by^2} = 0 \quad \int dy y^2 e^{-by^2} = \frac{1}{2b} \sqrt{\frac{\pi}{b}} \right\}$$

$$\Rightarrow U(x_0, x_0; T) = \left(\frac{1}{c} \sqrt{\frac{2\pi\hbar\epsilon}{-im}} \right) \left[1 - \frac{i\epsilon}{\hbar} V(x_0) + \frac{i\epsilon\hbar}{2m} \frac{\partial^2}{\partial x_0^2} + \mathcal{O}(\epsilon^2) \right] \\ \times U(x_0, x_0; T-\epsilon)$$

As $\epsilon \rightarrow 0$ makes sense wrt if $(\cdot) = 1$

$$\Rightarrow \mathcal{F}(\epsilon) = \sqrt{\frac{2\pi\hbar\epsilon}{-im}}$$

Now we can further expand $U(x_1, x_2; T - \epsilon)$:

$$U(x_1, x_2; T) = (1) [1 - \epsilon V(x_1) + \epsilon \bar{\partial}_2^2 + O(\epsilon^2)]$$

$$\alpha [1 - \epsilon \bar{\partial}_2^2 + O(\epsilon^2)] U(x_1, x_2; T)$$

match it to $O(\epsilon)$:

$$\epsilon \bar{\partial}_2 U(x_1, x_2; T) = [-\bar{\partial}_2^2 + V(x_1)] U(x_1, x_2; T)$$

$$= H U(x_1, x_2; T)$$

Schrodinger Equation !!

{ of course $U(x_1, x_2; T)$ is covered formula also }
 solutions has equation !!

What about initial condition ??

converged: As $T \rightarrow 0$ $\langle x_1 | e^{-iHT/\epsilon} | x_2 \rangle \rightarrow \delta(x_1 - x_2)$

path integral: As $T \rightarrow 0$ $\int \exp\left[\frac{i}{\epsilon} \int_0^T m(\dot{x}_1 - \dot{x}_2)^2 dt\right] \rightarrow \delta(x_1 - x_2)$

=> SAME LIMIT !!
 CONVERGENCE !!
 1 km scale!